

Now consider the Neumann eigenvalue problem for the  $p$ -Laplacian

$$-\Delta_p u = \Lambda_p^p |u|^{p-2} u \quad \text{in } \Omega$$

$$|\nabla u|^{p-2} \frac{\partial u}{\partial n} = 0 \quad \text{on } \partial\Omega,$$

with (first nontrivial eigenvalue)  $\Lambda_p^p$  minimizing

$$\mathcal{R}_p(v) := \frac{\int_{\Omega} |\nabla v|^p}{\int_{\Omega} |v|^p} \text{ on } W^{1,p}(\Omega) \cap \left\{ \int_{\Omega} |v|^{p-2} v = 0 \right\}. \quad (2)$$

What is known about  $\Lambda_p$  and what happens as  $p \rightarrow 1$  or  $p \rightarrow \infty$  ?

As  $p \rightarrow 1$ , formally we minimize  $\int_{\Omega} |Dv| dx$   
on functions satisfying  $\|v\|_{L^1(\Omega)} = 1$  and  $\int_{\Omega} \text{sign} v dx = 0$ .  
Minimizers can be chosen piecewise constant:  $c > 0$  on  $\Omega^+$  and  $-c$  on  $\Omega^-$ .

This constitutes a geometric partitioning problem:  
Divide  $\Omega$  into two disjoint subsets  $\Omega^+$  and  $\Omega^-$  of equal volume  
such that their relative perimeter  $P(\partial\Omega^+; \Omega)$  becomes minimal.

Problems of this nature are relevant in developing fast algorithms  
for parallel computing (Gajewski, Gärtner: Domain separation  
by means of sign changing eigenfunctions of p-Laplacians. 1998-2001)

For  $p \rightarrow \infty$  we have  $\Lambda_p \rightarrow \Lambda_\infty := 2/\text{diam}(\Omega)$ , while  $\lambda_\infty = 1/R(\Omega)$ .  
 The first nontrivial eigenfunctions  $u_p$  converge to a viscosity solution  $u_\infty$  of

$$\begin{cases} \min\{|\nabla u| - \Lambda u, -\Delta_\infty u\} = 0 & \text{in } \{u > 0\} \cap \Omega \\ \max\{-|\nabla u| - \Lambda u, -\Delta_\infty u\} = 0 & \text{in } \{u < 0\} \cap \Omega \\ -\Delta_\infty u = 0 & \text{in } \{u = 0\} \cap \Omega \\ \partial u / \partial \nu = 0 & \text{on } \partial\Omega. \end{cases} \quad (3)$$

While (3) was already derived in (Rossi, Saintier 14), we also prove that  $\Lambda \geq \Lambda_\infty$  for any  $\Lambda$  having a nontrivial solution to (3) and any convex  $\Omega$ . This shows that (for convex  $\Omega$ )  $\Lambda_\infty$  is in fact the first positive eigenvalue of (3).

The fact that  $\Lambda_\infty = 2/\text{diam}(\Omega)$  has interesting consequences.

**Corollary 1:** *If  $\Omega^*$  denotes a ball of same volume as  $\Omega$ , then the Szegő-Weinberger inequality  $\Lambda_\infty(\Omega) \leq \Lambda_\infty(\Omega^*)$  holds.*

For general  $p$  we do not know that  $\Lambda_p(\Omega) \leq \Lambda_p(\Omega^*)$ , except

a) for  $p = 2$  and general  $\Omega$  or

b) for  $p = \infty$  and convex  $\Omega$  (Cor.1).

c) For  $p = 1$  and convex plane  $\Omega$  this was an old conjecture of Polya  
(Esposito, Ferone, Kawohl, Nitsch, Trombetti 12).

**Corollary 2:** *For convex  $\Omega$  we have  $\Lambda_\infty(\Omega) \leq \lambda_\infty(\Omega)$ .*

*Moreover, equality holds only if  $\Omega$  is a ball.*

A stronger statement is known for finite  $p$ , namely  $\Lambda_p(\Omega) < \lambda_p(\Omega)$

a) for  $p = 2$  & general  $\Omega$  (Faber-Krahn & Cor. 1) or  $p = 1$  & convex plane  $\Omega$

b) for  $p \in (1, \infty)$  and convex  $\Omega$  (Brasco, Nitsch, Trombetti 14)

**Corollary 3:** *For convex  $\Omega$  no Neumann eigenfunction associated to  $\Lambda_\infty$  can have a closed nodal line inside  $\Omega$ .*

Else there is a nodal domain  $\Omega' \subset\subset \Omega$  so that

$\Lambda_\infty(\Omega) = \lambda_\infty(\Omega') = 1/R(\Omega') > 1/R(\Omega) = \lambda_\infty(\Omega)$  contradicting Cor. 2.

**Lemma 1:** For  $\Omega$  simply connected and Lipschitz  $\Lambda_p \rightarrow \Lambda_\infty := 2/\text{diam}(\Omega)$ , where  $\text{diam}$  denotes intrinsic diameter.

For  $x, y \in \Omega$  the intrinsic distance  $d_\Omega(x, y)$  is the length of the geodesic in  $\Omega$  connecting  $x$  to  $y$  and  $\text{diam}(\Omega) = \sup_{x, y \in \Omega} d_\Omega(x, y)$ .

Proof in 2 steps.

**Step 1:**  $\limsup_{p \rightarrow \infty} \Lambda_p \leq \Lambda_\infty$ .

Pick  $x_0 \in \Omega$  and adjust  $c_p \in \mathbb{R}$  so that  $w(x) = d_\Omega(x, x_0) - c_p$  is admissible test function for  $\mathcal{R}_p$ . Then

$$\Lambda_p \leq \mathcal{R}_p(w) = \left( \frac{1}{|\Omega|} \int_\Omega |d_\Omega(x, x_0) - c_p|^p \right)^{-1/p}.$$

Since  $0 \leq c_p \leq \text{diam}(\Omega)$ , up to subsequence  $c_p \rightarrow c_\infty$  and

$$\liminf_{p \rightarrow \infty} \left( \frac{1}{|\Omega|} \int_\Omega |d_\Omega(x, x_0) - c_p|^p \right)^{1/p} = \sup_{x \in \Omega} |d_\Omega(x, x_0) - c_\infty| \geq \text{diam}(\Omega)/2$$

**Step 2:**  $\liminf_{p \rightarrow \infty} \Lambda_p \geq \Lambda_\infty$ .

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For  $p > m > n$  the family of eigenfunctions  $u_p$  is uniformly bounded in  $W^{1,m}(\Omega)$  and equicont. so it converges in  $C^0$  and weakly in  $W^{1,m}$  to a limit  $u_\infty$ .

$$\frac{\|\nabla u_\infty\|_m}{\|u_\infty\|_m} \leq \liminf_{p \rightarrow \infty} \frac{\|\nabla u_p\|_m}{\|u_p\|_m} \leq \liminf_{p \rightarrow \infty} \frac{\|\nabla u_p\|_p}{\|u_p\|_m} = \liminf_{p \rightarrow \infty} \Lambda_p \frac{\|u_p\|_p}{\|u_p\|_m}$$

Sending  $m \rightarrow \infty$  gives  $\liminf_{p \rightarrow \infty} \Lambda_p \geq \frac{\|\nabla u_\infty\|_\infty}{\|u_\infty\|_\infty}$   
so it remains to estimate  $\|u_\infty\|_\infty$  in terms of  $\|\nabla u_\infty\|_\infty$ .

Condition  $\int_\Omega |u_p|^{p-2} u_p = 0$  implies  $\sup u_\infty = -\inf u_\infty$

$$\|u_\infty\|_\infty = \frac{1}{2}(\sup u_\infty - \inf u_\infty) \leq \frac{1}{2} \text{diam}(\Omega) \|\nabla u_\infty\|_\infty.$$

Thus

$$\liminf_{p \rightarrow \infty} \Lambda_p \geq \frac{\|\nabla u_\infty\|_\infty}{\|u_\infty\|_\infty} \geq \frac{2}{\text{diam}(\Omega)} = \Lambda_\infty$$



## Observation:

Suppose  $\|u_\infty\|_\infty$  is scaled to 1,  $u(x) = \inf u$  and  $u(y) = \sup u$ .

Then  $u_\infty$  increases with constant slope  $\Lambda_\infty$  along the geodesic from  $x$  to  $y$ .

In particular in the one-dimensional case that  $\Omega = (-1, 1)$  we have  $\Lambda_\infty = 1$ , and the function  $u(x) = x$  is a viscosity solution of (3).

$$\begin{cases} \min\{|\nabla u| - \Lambda u, -\Delta_\infty u\} = 0 & \text{in } \{u > 0\} \cap \Omega \\ \max\{-|\nabla u| - \Lambda u, -\Delta_\infty u\} = 0 & \text{in } \{u < 0\} \cap \Omega \\ -\Delta_\infty u = 0 & \text{in } \{u = 0\} \cap \Omega \\ \partial u / \partial \nu = 0 & \text{on } \partial\Omega. \end{cases}$$

While it solves the differential equation  $-\Delta_\infty u = u'' = 0$  in the classical (and thus in the viscosity) sense, it clearly violates the boundary condition  $u'(\pm 1) = 0$  in the classical sense.

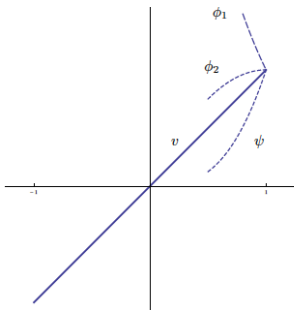
$u(x) = x$  clearly violates the boundary condition  $u'(\pm 1) = 0$  in the classical sense. To check the Neumann condition in the right end point  $x = 1$  in the viscosity sense, one must verify

$$\min\{\min\{|\phi'| - \Lambda\phi, -|\phi'|^2\phi''\}, \phi'\}(1) \leq 0 \quad (4)$$

for any  $C^2$  test function  $\phi$  touching  $u$  in  $x = 1$  from above, and

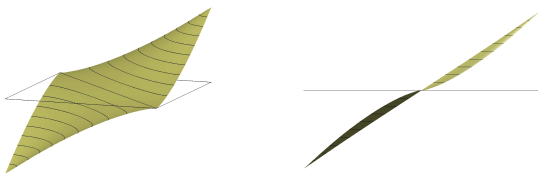
$$\max\{\min\{|\psi'| - \Lambda\psi, -|\psi'|^2\psi''\}, \psi'\}(1) \geq 0 \quad (5)$$

for any smooth test function  $\psi$  touching  $u$  from below.

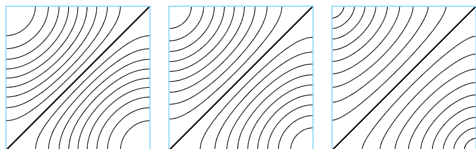




Moreover, on a rectangle  $u_\infty$  has constant (and maximal) slope along a diagonal.



Numerical simulation of  $u_{15}$  and side view in diagonal direction.



Level plots of  $u_p$  for  $p = 2$ ,  $p = 4$  and  $p = 8$

For a square and  $p = 2$  (the linear case)  $\Lambda_2$  is a double eigenvalue and there are solutions with diagonal and horizontal or vertical nodal lines.

There are one-dimensional solutions also for  $p = \infty$  but they are associated to higher eigenvalues than  $\Lambda_\infty$ .

Take  $\Omega = (-1, 1)^2$  and  $u(x) = x_1$ . Then  $u \in C^2(\Omega)$ ,  $-\Delta_\infty u = 0$  in  $\Omega$ .

Now PDE in (3) is satisfied if also  $1 = |\nabla u| \geq \Lambda u$  on  $\{u > 0\}$ , that implies

$$\Lambda \leq 1.$$

The Neumann bc. is satisfied in classical sense on horizontal parts of  $\partial\Omega$ . However, for Neumann bc. to hold in the viscosity sense on the right part, we must verify

$$\min\{\min\{|\nabla\phi| - \Lambda\phi, -\Delta_\infty\phi\}, \partial\phi/\partial\nu\}(x_0) \leq 0$$

for any  $C^2$  test function  $\phi$  touching  $u$  in  $x_0 \in \partial\Omega$  from above, and

$$\max\{\min\{|\nabla\psi| - \Lambda\psi, -\Delta_\infty\psi\}, \partial\psi/\partial\nu\}(x_0) \geq 0$$

for any smooth test function  $\psi$  touching  $u$  from below.

Recall  $|\nabla u| = \partial u/\partial\nu = 1$  at  $x_1 = 1$ . Therefore only the very first constraint for  $\psi$  is active and implies

$$\Lambda \geq 1.$$

This shows that  $u(x) = x_1$  is a viscosity solution to (3) with eigenvalue  $\Lambda = 1$ , but

$$\Lambda = 1 > \frac{1}{\sqrt{2}} = \frac{2}{\text{diam}(\Omega)} = \Lambda_\infty.$$

# Eigenvalue problem for the normalized $p$ -Laplacian

Also the eigenvalue problem

$$-\Delta_p^N u = \lambda_p u \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega \quad (6)$$

has been studied for  $p \in (1, \infty)$  by Birindelli/Demengel, Rossi et al and for  $p = \infty$  by Juutinen. Now one has the estimate

$$\left(\frac{\pi}{2\rho(\Omega)}\right)^2 \leq \lambda_\infty(\Omega) \leq \left(\frac{\pi}{2R(\Omega)}\right)^2,$$

in which  $\rho(\Omega)$  and  $R(\Omega)$  denote the outer and inradius of  $\Omega$ , and which becomes sharp for a ball of radius  $R$ .

For  $p \geq 2$  the first eigenfunction is again unique modulo scaling. At least for starshaped domains  $\Omega$  there is convergence of the first eigenvalue  $\lambda_p$  to  $\lambda_\infty$  as  $p \rightarrow \infty$ , but the limit  $p \rightarrow 1$  appears to be open.

### Open Problem

Let  $\Omega = B_R(0)$  be a ball and  $u$  a radially symmetric eigenfunction. One should also expect many nonradial eigenfunctions to exist, but for  $p \neq 2$  we are not aware of any results in this direction, not even in two dimensions. For radial functions and with the Ansatz  $u(x) = v(|x|)$  the eigenvalue problem (6) transforms into

$$(p-1)v''(r) + \frac{n-1}{r}v'(r) + p\lambda v(r) = 0 \quad \text{in } (0, R), \quad v'(0) = 0 = v(R). \quad (7)$$

The formal limit  $p \rightarrow 1$  in (7) is

$$\frac{n-1}{r}v'(r) + \lambda v(r) = 0 \quad \text{in } (0, R), \quad v'(0) = 0 = v(R). \quad (8)$$

with the viscosity solution  $v(r) = v(0) e^{-\frac{\lambda}{2(n-1)}r^2}$ , which violates the boundary condition in the classical sense at  $r = R$ , but since it solves (8) there in the classical sense, it still satisfies the boundary condition  $v(R) = 0$  in the viscosity sense. It is therefore reasonable to expect that also for more general domains the first eigenfunctions in (32) converge (as  $p \rightarrow 1$ ) to a viscosity solution of

$$(n-1)Hv_\nu + \lambda v = 0 \quad \text{in } \Omega, \quad v = 0, \quad \text{on } \partial\Omega, \quad (9)$$

$$(p-1)v''(r) + \frac{n-1}{r}v'(r) + p\lambda v(r) = 0 \quad \text{in } (0, R), \quad v'(0) = 0 = v(R). \quad (7)$$

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and that they develop a boundary layer.

For  $p > 1$  problem

$$(p-1)v''(r) + \frac{n-1}{r}v'(r) + p\lambda v(r) = 0 \quad \text{in } (0, R), \quad v'(0) = 0 = v(R). \quad (7)$$

is a Bessel type equation, and this observation was used by Kawohl/Krömer/Kurtz to explicitly derive a countable and complete orthonormal system of eigenfunctions to (7) in a suitably weighted  $L^2$ -space. These form a countable, but presumably incomplete sequence of eigenfunctions to (6)

$$-\Delta_p^N u = \lambda_p u \quad \text{in } B_R(0), \quad u = 0 \quad \text{on } \partial B_R(0).$$



Other open issues concerning  $p$ -Laplacian:

- Uniqueness of first Dirichlet eigenfunction for  $p = 1$  or  $\infty$  and convex  $\Omega$ ?  
There is nonuniqueness for nonconvex  $\Omega$ .
- Nodal pattern on disc for any  $p \neq 2$ ?  
Not radial for both second Dirichlet and Neumann eigenf., but diameter?
- Hot spot conjecture true for general  $p \in (1, \infty)$  and convex  $\Omega$ ?  
Known for  $p = 1$  and  $\infty$ . Open even for  $p = 2$ , see Polymath thread of T.Tao.